



# A Three-Point Boundary Value Problem with an Integral Condition for Parabolic Equations with the Bessel Operator

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**Abstract**—In this paper, we study a three-point boundary value problem with an integral condition for a class of parabolic equation with Bessel operator. The existence and uniqueness of the solution in functional weighted Sobolev space are proved. The proof is based on two sided *a priori* estimates and the density of the range of the operator generated by the considered problem. © 2000 Elsevier Science Ltd. All rights reserved.

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## 1. INTRODUCTION

In the rectangle  $\Omega = (0, T) \times (0, \ell)$ , we consider the equation

$$\mathcal{L}u = \frac{\partial u}{\partial t} - \left( \frac{\partial^2 u}{\partial x^2} + \frac{1}{x} \frac{\partial u}{\partial x} \right) = f(t, x). \quad (1)$$

To equation (1), we add the initial condition,

$$lu = u(0, x) = \varphi(x), \quad x \in (0, \ell), \quad (2)$$

the Dirichlet condition,

$$u(t, \ell) = 0, \quad t \in (0, T), \quad (3)$$

and integral condition,

$$\int_{\ell_1}^{\ell} u(t, \xi) d\xi = 0, \quad 0 \leq \ell_1 \leq \ell, \quad t \in (0, T). \quad (4)$$

Here, we assumed that the function  $\varphi$  satisfies the conditions given in (3) and (4), i.e.,

$$\varphi(\ell) = 0, \quad \int_{\ell_1}^{\ell} \varphi(x) dx = 0.$$

The importance of problems with integral conditions has been pointed out by Samarskii [1]. Regular case of this problem is studied in [2]. The problem where the equation contain an operator of the form  $\frac{\partial}{\partial x}(a(t, x)\frac{\partial u}{\partial x})$ , instead of Bessel operator, is treated in [3]. Similar problems for second-order parabolic equations are investigated by the potential method in [4]. Two-point boundary value problems for parabolic equations, with an integral condition, are investigated using the energy inequalities method in [5–9] and the Fourier method [10].

We associate to problems (1)–(4) the operator  $L = (\mathcal{L}, l)$ , defined from  $E$  into  $F$ , where  $E$  is the Banach space of functions  $u \in L_2(\Omega)$ , satisfying (3) and (4), with the finite norm

$$\|u\|_E^2 = \int_{\Omega} \theta(x) \left| \frac{\partial u}{\partial t} \right|^2 dx dt + \int_{\Omega} \left| \frac{\partial}{\partial x} \left( x \frac{\partial u}{\partial x} \right) \right|^2 dx dt + \sup_{0 \leq t \leq T} \int_0^{\ell} \theta(x) \left| \frac{\partial u}{\partial x} \right|^2 dx, \quad (5)$$

and  $F$  is the Hilbert space of vector-valued functions  $\mathcal{F} = (f, \varphi)$  obtained by completion of the space  $L_2(\Omega) \times W_2^2(0, \ell)$  with respect to the norm

$$\|\mathcal{F}\|_F^2 = \|(f, \varphi)\|_F^2 = \int_{\Omega} \theta(x) |f(t, x)|^2 dx dt + \int_0^{\ell} \theta(x) \left| \frac{d\varphi}{dx} \right|^2 dx, \quad (6)$$

where

$$\theta(x) = \begin{cases} \ell_1 x, & 0 < x \leq \ell_1, \\ x^2, & \ell_1 \leq x < \ell. \end{cases} \quad (7)$$

Using the energy inequalities method proposed in [9], we establish two-sided *a priori* estimates. Then we prove that the operator  $L$  is a linear homeomorphism between the spaces  $E$  and  $F$ .

## 2. TWO-SIDED A PRIORI ESTIMATES

THEOREM 1. *The following a priori estimate:*

$$\|Lu\|_F \leq C \|u\|_E, \quad (8)$$

holds for any function  $u \in E$ , where the constant  $C$  is independent of  $u$ .

PROOF. Using equation (1) and initial condition (2), we obtain

$$\int_{\Omega} \theta(x) |\mathcal{L}u|^2 dx \leq 2 \int_{\Omega} \left[ \theta(x) \left| \frac{\partial u}{\partial t} \right|^2 + \left| \frac{\partial}{\partial x} \left( x \frac{\partial u}{\partial x} \right) \right|^2 \right] dx dt \quad (9)$$

and

$$\int_0^{\ell} \theta(x) \left| \frac{dlu}{dx} \right|^2 dx \leq \sup_{0 \leq t \leq T} \int_0^{\ell} \theta(x) \left| \frac{\partial u}{\partial x} \right|^2 dx, \quad (10)$$

combining inequalities (9) and (10), we obtain (8) for  $u \in E$ .

THEOREM 2. *For any function  $u \in E$ , we have the inequality*

$$\|u\|_E \leq C \|Lu\|_F, \quad (11)$$

where the constant  $C$  is independent of  $u$ .

PROOF. We denote  $D(L) = \{u \in E : x \frac{\partial^2 u}{\partial x \partial t} \in L_2(\Omega)\}$ ,

$$J_x g = \int_x^{\ell} g(t, \xi) d\xi,$$

and

$$Mu = \begin{cases} \ell_1 x u, & 0 < x \leq \ell_1, \\ x^2 u + x J_x u, & \ell_1 \leq x < \ell, \end{cases}$$

multiplying equation (1) by  $M \overline{\frac{\partial u}{\partial t}}$  and integrating over  $\Omega^\tau$ , where  $\Omega^\tau = (0, \tau) \times (0, \ell)$ , we get

$$\begin{aligned} \operatorname{Re} \int_0^\tau \int_0^\ell \mathcal{L}u M \overline{\frac{\partial u}{\partial t}} dx dt &= \int_0^\tau \int_0^\ell \theta(x) \left| \frac{\partial u}{\partial t} \right|^2 dx dt \\ &+ \operatorname{Re} \int_0^\tau \int_0^\ell \theta(x) \frac{\partial u}{\partial x} \frac{\partial^2 u}{\partial x \partial t} dx dt + \frac{1}{2} \int_0^\tau \int_{\ell_1}^\ell \left| J_x \frac{\partial u}{\partial t} \right|^2 dx dt. \end{aligned} \quad (12)$$

Integrating by parts  $\operatorname{Re} \int_0^\tau \int_0^\ell \theta(x) \frac{\partial u}{\partial x} \frac{\partial^2 u}{\partial x \partial t} dx dt$ , we obtain

$$\begin{aligned} \operatorname{Re} \int_0^\tau \int_0^\ell \mathcal{L}u M \overline{\frac{\partial u}{\partial t}} dx dt &+ \frac{1}{2} \int_0^\ell \theta(x) \left| \frac{du}{dx} \right|^2 dx = \int_0^\tau \int_0^\ell \theta(x) \left| \frac{\partial u}{\partial t} \right|^2 dx dt \\ &+ \frac{1}{2} \int_0^\ell \theta(x) \left| \frac{\partial u(\tau, x)}{\partial x} \right|^2 dx + \frac{1}{2} \int_0^\tau \int_{\ell_1}^\ell \left| J_x \frac{\partial u}{\partial t} \right|^2 dx dt, \end{aligned} \quad (13)$$

and using the elementary inequalities, we get

$$\begin{aligned} \operatorname{Re} \int_0^\tau \int_0^\ell \mathcal{L}u M \overline{\frac{\partial u}{\partial t}} dx dt &\leq \frac{3}{2} \int_0^\tau \int_0^\ell \theta(x) |\mathcal{L}u|^2 dx dt \\ &+ \frac{1}{4} \int_0^\tau \int_0^\ell \theta(x) \left| \frac{\partial u}{\partial t} \right|^2 dx dt + \frac{1}{2} \int_0^\tau \int_{\ell_1}^\ell \left| J_x \frac{\partial u}{\partial t} \right|^2 dx dt. \end{aligned} \quad (14)$$

From equation (1), we have

$$\frac{1}{4} \int_0^\tau \int_0^\ell \left| \frac{\partial}{\partial x} \left( x \frac{\partial u}{\partial x} \right) \right|^2 dx dt \leq \frac{1}{2} \int_0^\tau \int_0^\ell \theta(x) \left| \frac{\partial u}{\partial t} \right|^2 dx dt + \frac{1}{2} \int_0^\tau \int_0^\ell \theta(x) |\mathcal{L}u|^2 dx dt. \quad (15)$$

Combining inequalities (14) and (15), we get

$$\begin{aligned} \frac{1}{4} \int_0^\tau \int_0^\ell \theta(x) \left| \frac{\partial u}{\partial t} \right|^2 dx dt &+ \frac{1}{2} \int_0^\ell \theta(x) \left| \frac{\partial u(\tau, x)}{\partial x} \right|^2 dx \\ &+ \frac{1}{4} \int_0^\tau \int_0^\ell \left| \frac{\partial}{\partial x} \left( x \frac{\partial u}{\partial x} \right) \right|^2 dx dt \leq 2 \int_0^\tau \int_0^\ell \theta(x) |\mathcal{L}u|^2 dx dt + \frac{1}{2} \int_0^\ell \theta(x) \left| \frac{du}{dx} \right|^2 dx. \end{aligned} \quad (16)$$

The right-hand side of (16) is independent of  $\tau$ , hence replacing the left-hand side by its upper bound with respect to  $\tau$ , in the interval  $[0, T]$ , we obtain the desired inequality.

### 3. SOLVABILITY OF THE PROBLEM

From estimates (8) and (11), it follows that the operator  $L : E \mapsto F$  is continuous and its range is closed in  $F$ . Therefore, the inverse operator  $L^{-1}$  exists and continuous, from the closed subspace  $R(L)$  onto  $E$ , which means, that  $L$  is a homeomorphism from  $E$  to  $R(L)$  and so to have a unique solution to the problem it remains to show that  $R(L) = F$ .

The proof is based on the following theorem.

**THEOREM 3.** Let  $D_0(L) = \{u \in D(L) / lu = 0\}$ . If

$$\int_\Omega \theta(x) \mathcal{L}u \bar{\omega} dx dt = 0, \quad (17)$$

for  $\omega$  such that  $x\omega \in L_2(\Omega)$  and all  $u \in D_0(L)$ , then  $\omega$  vanishes almost everywhere on  $\Omega$ .

PROOF. Putting into (17)  $h = \frac{\partial u}{\partial t}$ ; where  $h, x \frac{\partial h}{\partial x}, \frac{\partial}{\partial x}(x \frac{\partial h}{\partial x}) \in L_2(\Omega)$ , and  $h$  verifies the boundary conditions (3) and (4), we get

$$\begin{aligned} \int_{\Omega} \theta(x) h \bar{\omega} dx dt &= \int_{\Omega} \frac{1}{x} \theta(x) \frac{\partial}{\partial x} \left( x \frac{\partial h}{\partial x} \right) J_t \bar{\omega} dx dt = \int_{\Omega} \frac{1}{x^2} \theta(x) \frac{\partial}{\partial x} \left( x \frac{\partial}{\partial x} (xh) \right) J_t \bar{\omega} dx dt \\ &\quad - 2 \int_{\Omega} \frac{1}{x^2} \theta(x) \frac{\partial}{\partial x} (xh) J_t \bar{\omega} dx dt + \int_{\Omega} \frac{1}{x^3} \theta(x) (xh) J_t \bar{\omega} dx dt, \end{aligned} \quad (18)$$

where  $J_t \omega = \int_t^T \omega(\tau, x) d\tau$ .

The left-hand side of (18) shows that the mapping

$$xh \longmapsto \int_{\Omega} \frac{1}{x} \theta(x) \frac{\partial}{\partial x} \left( x \frac{\partial h}{\partial x} \right) J_t \omega dx dt$$

is a continuous linear functional. From the right-hand side of (18), there follows that this is true if the function  $\omega$  has the following properties:

$$\frac{\theta(x)}{x^3} J_t \omega, \frac{\theta(x)}{x^2} \frac{\partial}{\partial x} (J_t \omega), \frac{\theta(x)}{x^2} \frac{\partial}{\partial x} \left( x \frac{\partial}{\partial x} (J_t \omega) \right) \in L_2(\Omega) \quad \text{and} \quad J_t \omega|_{x=\ell} = 0,$$

In terms of the given function  $\omega$ , from equality (17), we define the function

$$v(t, x) = \begin{cases} J_t \omega, & 0 < x \leq \ell_1, \\ - \int_x^{\ell} J_t \omega(t, \xi) \frac{1}{\xi} d\xi + J_t \omega, & \ell_1 \leq x < \ell. \end{cases}$$

Hence, it follows that  $\int_x^{\ell} v(t, \xi) d\xi = - \int_x^{\ell} d\xi \int_{\xi}^{\ell} J_t \omega(t, \eta) d\eta + \int_x^{\ell} J_t \omega(t, \xi) d\xi = x \int_x^{\ell} J_t \omega(t, \xi) (1/\xi) d\xi$ ,

$$\theta(x) J_t \omega = Mv \quad \text{and} \quad \int_{\ell_1}^{\ell} v(t, x) dx = 0, \quad (19)$$

replacing  $h = J_t^* v$  in (18), where  $J_t^* v = \int_0^t v(\tau, x) d\tau$ , we get

$$\int_{\Omega} J_t^* v \frac{\partial}{\partial t} (M\bar{v}) dx dt = - \int_{\Omega} \frac{1}{x} \frac{\partial}{\partial x} \left( x \frac{\partial}{\partial x} J_t^* v \right) M\bar{v} dx dt, \quad (20)$$

doing the same calculations as in Theorem 2, where replacing  $\frac{\partial u}{\partial t}$  by  $v$ , and taking into account (19), we obtain

$$\int_{\Omega} \theta(x) |v|^2 dx dt + \frac{1}{2} \int_0^T \int_{\ell_1}^{\ell} |J_x v|^2 dx dt \leq 0,$$

then  $v = 0$ , hence  $\omega = 0$ . This proves Theorem 3.

**THEOREM 4.** *The range  $R(L)$  of the operator  $L$  coincides with  $F$ .*

PROOF. Since  $F$  is a Hilbert space, we have  $R(L) = F$  if and only if the equality

$$\int_{\Omega} \theta(x) \mathcal{L} u \bar{f} dx dt + \int_0^{\ell} \left( \theta(x) \frac{dlu}{dx} \frac{d\bar{\varphi}}{dx} + lu \bar{\varphi} \right) dx = 0, \quad (21)$$

which gives that  $\mathcal{F} = (f, \varphi) = 0$ . In particular, if we put  $lu = 0$  in (21), then we conclude by Theorem 3 that  $f = 0$ , so it follows from (21)  $\int_0^{\ell} \theta(x) \frac{dlu}{dx} \frac{d\bar{\varphi}}{dx} + lu \bar{\varphi} dx = 0$ . Since the range of the trace operator  $l$  is dense in the Hilbert space with the norm  $[\int_0^{\ell} (\theta(x) |\frac{d\varphi}{dx}|^2 + |\varphi|^2) dx]^{1/2}$ , hence  $\varphi = 0$ . Consequently,  $\mathcal{F} = 0$ .

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